## M6a: Open Channel Flow <br> (Manning's Equation, Partially Flowing Pipes, and Specific Energy)

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Chin 2006; Figure 3.1

Steady, Non-Uniform Flow in an Open Channel


The momentum equation can be used to derive the expression for $x$ shear stress:

$$
\tau_{o}=\gamma R S_{f}
$$

$\mathrm{y}=$ specific weight of water ( $62.4 \mathrm{lbs} / \mathrm{tt} 3$ )
$\mathrm{R}=$ hydraulic radius (ft)
$\mathrm{Sf}=$ hydraulic slope (ft/ft) (slope of the energy grade line, or the friction slope)
Chin 2006; Figure 3.2

SI units (m/s; m)

$$
V=\frac{1}{n} R^{2 / 3} S_{f}^{1 / 2}
$$

U.S. Customary units (ft/sec; ft)

$$
V=\frac{1.486}{n} R^{2 / 3} S_{f}{ }^{1 / 2}
$$

Where Sf is the slope of the energy grade line (friction slope). If the channel slope is used, then implies uniform flow, which is rare.
Only valid for hydraulically rough flow, when: $n^{6} \sqrt{R S_{f}} \geq 1.9 \times 10^{-13}$ (SI units)

## Manning's Equation

## Example:

- Determine the flow rate in a rectangular concrete channel with a width of 3 m and a HGL slope of $0.001 \mathrm{~m} / \mathrm{m}$ when the depth of flow is 1.5 m . Assume $\mathrm{n}=0.014$.


Given: $\quad \mathrm{n}=0.014$ (concrete channel)
$\mathrm{L}=3 \mathrm{~m}$ (width of channel)
$\mathrm{w}=1.5 \mathrm{~m}$ (depth of flow)
$\mathrm{s}_{\mathrm{f}}=0.001 \mathrm{~m} / \mathrm{m}$

## Manning's Equation

- Use the Manning's equation:

$$
Q=\frac{A}{n} R^{2 / 3} S_{f}^{1 / 2}
$$

Need A (cross-sectional area of flow):
A = Lw
Substituting:

$$
\begin{aligned}
& A=(3 \mathrm{~m})(1.5 \mathrm{~m}) \\
& A=4.5 \mathrm{~m}^{2}
\end{aligned}
$$

## Manning's Equation

- Need R (hydraulic radius):

$$
\mathrm{R}=\mathrm{A} / \mathrm{P}
$$

- Need P, the wetted perimeter (noted on drawing by thicker lines).

$$
P=L+2 w
$$

$$
\begin{aligned}
& P=(3 \mathrm{~m})+2(1.5 \mathrm{~m}) \\
& P=6 \mathrm{~m}
\end{aligned}
$$

Substituting into equation for hydraulic radius:

$$
\begin{aligned}
& \mathrm{R}=\mathrm{A} / \mathrm{P} \\
& \mathrm{R}=\left(4.5 \mathrm{~m}^{2}\right) /(6 \mathrm{~m}) \\
& \mathrm{R}=0.75 \mathrm{~m}
\end{aligned}
$$

## Manning's Equation

- Substituting into Manning's equation:
Substituting:

$$
\begin{aligned}
& Q=\frac{\left(4.5 m^{2}\right)(0.75 m)^{2 / 3}(0.001 \mathrm{~m} / \mathrm{m})^{1 / 2}}{0.014} \\
& Q=8.39 \mathrm{~m}^{3} / \mathrm{sec}
\end{aligned}
$$

## Manning's Equation

## Example:

- Given a V-shaped channel with a HGL slope of 0.001, a top width of 10 feet, and a depth of 5 feet, determine the velocity of flow using the Manning's equation. Find the discharge in both $\mathrm{ft}^{3} / \mathrm{sec}$ (cfs) and $\mathrm{m}^{3} / \mathrm{sec}$ (cms).



## Manning's Equation

Substituting:

$$
\begin{aligned}
& \mathrm{R}=\left(25 \mathrm{ft}^{2}\right) / 14.14 \mathrm{ft} \\
& \mathrm{R}=1.77 \mathrm{ft}
\end{aligned}
$$

Assume that the channel is concrete-lined with a Manning's n of 0.015 .

$$
\begin{gathered}
Q=\frac{1.49}{n} A R^{2 / 3} S^{1 / 2} \\
Q=\frac{1.49}{0.015}\left(25 \mathrm{ft}^{2}\right)(1.77 \mathrm{ft})^{2 / 3}(0.001)^{1 / 2}
\end{gathered}
$$

$$
Q=114.9 \mathrm{cfs}=115 \mathrm{ft}^{3} / \mathrm{sec}(0.3048 \mathrm{~m} / \mathrm{ft})^{3}=32.5
$$

$$
\mathrm{m}^{3} / \mathrm{sec}
$$

## Manning's Equation

## Example:

- Find the dimensions of a rectangular concrete channel to carry a flow of $150 \mathrm{~m}^{3} / \mathrm{sec}$, with a HGL slope of 0.015 and a mean velocity of $10.2 \mathrm{~m} / \mathrm{sec}$.


Given:
$\mathrm{Q}=150 \mathrm{~m}^{3} / \mathrm{sec}$
$V=10.2 \mathrm{~m} / \mathrm{sec}$
$\mathrm{S}_{\mathrm{f}}=0.015$
Assume: Manning's $\mathrm{n}=0.013$ (concrete channel)

## Manning's Equation

- Have everything needed to solve Manning's equation for the hydraulic radius, R :

$$
\begin{aligned}
& V=\frac{1}{n} R^{2 / 3} S^{1 / 2} \\
& R^{2 / 3}=\frac{V n}{S^{1 / 2}} \\
& R=\left(\frac{V n}{S^{1 / 2}}\right)^{3 / 2} \\
& R=\left(\frac{(10.2 m / \mathrm{sec})(0.013}{(0.015)^{1 / 2}}\right)^{3 / 2} \\
& R=1.13 m
\end{aligned}
$$

## Manning's Equation

By definition, $R=$ cross-sectional area of flow/wetted perimeter

$$
\mathrm{A}=(\text { Base of channel })(\text { Depth of flow })
$$

P = (Base of channel) +2 (Depth of flow)
Substituting:

$$
\mathrm{R}=\frac{(\text { Base })(\text { Depth })}{[\text { Base }+2 \text { Depth }]}=1.13 \mathrm{~m}
$$

By the Continuity Equation:

$$
\begin{aligned}
& A=\frac{Q}{V}=(\text { Base })(\text { Depth }) \\
& A=\frac{150 \mathrm{~m}^{3} / \mathrm{sec}}{10.2 \mathrm{~m} / \mathrm{sec}} \\
& A=14.7 \mathrm{~m}^{2}=(\text { Base })(\text { Depth })
\end{aligned}
$$

Figure 5-8 (Chow 1959) can be used to significantly shorten the calculation effort for the design of channels. This figure is used to calculate the normal depth ( y ) of a channel based on the channel side slopes and known flow and channel characteristics, using the Manning's equation in the following form:

$$
A R^{\frac{2}{3}}=\frac{n Q}{1.49 S^{0.5}}
$$

Initial channel characteristics that must be know include: z (the side slope), and b (the channel bottom width, assuming a trapezoid). It is easy to examine several different channel options ( $z$ and $b$ ) by calculating the normal depth (y) for a given peak discharge rate, channel slope, and roughness. The most practical channel can then be selected from the alternatives.

## Manning's Equation

- Have two equations and two unknowns:

$$
14.7 \mathrm{~m}^{2}=(\text { Base })(\text { Depth })
$$

$$
\begin{aligned}
& \text { Base }=\frac{14.7}{\text { Depth }} \\
& 1.13 m=\frac{14.7 m^{2}}{\text { Base }+2 \text { Depth }} \\
& \text { Base }+2 \text { Depth }=13.0 m \\
& \left(\frac{14.7}{\text { Depth }}\right)+2 \text { Depth }=13.0
\end{aligned}
$$

$$
14.7+2 \text { Depth }^{2}=13 \text { Depth }
$$

$$
2 \text { Depth }^{2}-13 \text { Depth }+14.7=0
$$

Two possible solutions to quadratic (both are correct):

| Base $=2.9 \mathrm{~m}$ | Depth of Flow $=5.04 \mathrm{~m}$ |
| :--- | :--- |
| Base $=10.1 \mathrm{~m}$ | Depth of Flow $=1.46 \mathrm{~m}$ |

Depth of Flow $=1.46 \mathrm{~m}$


## Composite Manning's n Estimate, Example 3.2 (Chin 2006)

A floodplain (next slide) can be divided into seven sections as shown below. Use the various formula in the table to estimate the composite roughness value for this channel.

|  |  | Formula | Reference |
| :---: | :---: | :---: | :---: |
|  |  | $n_{c}=\left(\frac{\sum_{i=1}^{N} P_{i} P_{i}^{3 / 2}}{\sum_{i=1}^{N} P_{i}}\right)^{2 / 3}$ | Horton (1933), Einstein (1934)* |
| Section | $n$ | $\left(\sum_{i=1}^{N} P_{i} n_{i}^{2}\right)^{1 / 2}$ |  |
| 1. | 0.040 | $\left(\sum_{i=1}^{N} P_{i}\right)^{1 / 2}$ | Muhlhofer (1933), Einstein and Banks (1951) |
| 2 | 0.030 | $P R^{3 / 3}$ |  |
| 3 | 0.015 |  | Loter (1933)' |
| 4 | 0.013 | $\sum_{i=1}^{n} \frac{P_{1} R_{i}}{n_{i}}$ |  |
| 5 | 0.017 |  |  |
| 6 | 0.035 | $\sum_{i=1}^{N} P_{i j} j^{1 / 2}$ | Krishnamurhy and Christensen (1972) |
| 7 | 0.060 | Formula assames that the mean flow in each of the subareas is cqual to the mean flow velocity. I $P$ and $R$ are the perimeter and lydraulic radius of the entire cross-section, reppectively. ${ }^{t} y$ is the werage flow depth in Section ; |  |

Floodplain showing seven separate sections corresponding to different values of $n$.


Figure 3.5, Chin 2006

These values are used with the prior equations to result in the following estimates for Manning's $\mathbf{n}$ :

| Formula | $\boldsymbol{n}_{\boldsymbol{e}}$ |
| :--- | :---: |
| Horton/Einstein | 0.033 |
| Muhlhofer/Einstein and Banks | $\mathbf{0 . 0 3 3}$ |
| Lotter | $\mathbf{0 . 0 2 2}$ |
| Krishnamurthy and Christensen | $\mathbf{0 . 0 2 6}$ |

The estimates of the composite $n$ values can vary considerably, resulting in similar differences in predicted discharges.

As shown earlier, the Manning's equation can also be also used to predict flows in pipes. Drainage systems are typically designed as open channel flows in circular pipes, although other cross-sectional shapes are used.

Charts or tables can be used to help predict the flow conditions in these systems when the pipes are not flowing full.

Sewers Flowing Partly Full

From: Metcalf and Eddy, Inc and George Tchobanoglous Wastewater Engineering: Collection and Pumping of umping of Wastewate McGraw-H Inc. 1981

## Sewers Flowing Partly Full

Example:

- Determine the depth of flow and velocity in a sewer with a diameter of 300 mm having a HGL slope of $0.005 \mathrm{~m} / \mathrm{m}$ with an n value of 0.015 when discharging $0.01 \mathrm{~m}^{3} / \mathrm{sec}$.

Given: $\quad D=300 \mathrm{~mm}$
$\mathrm{S}_{\mathrm{f}}=0.005 \mathrm{~m} / \mathrm{m}$
$\mathrm{n}=0.015$
$\mathrm{Q}=0.01 \mathrm{~m}^{3} / \mathrm{sec}$

## Sewers Flowing Partly Full

- Use the modified Manning's equation for partly full sewers:

$$
\begin{aligned}
& Q=\left(\frac{K^{\prime}}{n}\right) D^{8 / 3} S^{1 / 2} \\
& \text { Rearranging : } \\
& K^{\prime}=\frac{n Q}{D^{8 / 3} S^{1 / 2}} \\
& \text { Substituting : } \\
& K^{\prime}=\frac{(0.015)\left(0.01 \mathrm{~m}^{3} / \mathrm{sec}\right)}{(0.3 m)^{8 / 3}(0.005 \mathrm{~m} / \mathrm{m})^{1 / 2}} \\
& K^{\prime}=0.0526
\end{aligned}
$$

## Sewers Flowing Partly Full

- To calculate velocity at depth of water of 84 mm , need to use continuity equation:

$$
Q=V A
$$

- Using Manning's partial flow diagram (assuming a constant n ):

$$
\text { At } \mathrm{d} / \mathrm{D}=0.28
$$

## Sewers Flowing Partly Full

- Using Table 2-5 (equation in terms of diameter of pipe):

Close to $\mathrm{K}^{\prime}=0.0534$
Therefore, $\mathrm{d} / \mathrm{D}=0.28$

- Substituting:
$d /(0.3 \mathrm{~m})=0.28$
Depth of flow, $d=0.084 \mathrm{~m}$


Figure 2-16 Hydraulic elements for circular sewers [10].
From: Metcalf and Eddy, Inc. and George Tchobanoglous. Wastewater Engineering. Collection and Pumping of Wastewater. McGraw-Hill, Inc. 1981.

## Sewers Flowing Partly Full

- Using Manning's partial flow diagram (assuming a constant $n$ ):

$$
\text { At } \mathrm{d} / \mathrm{D}=0.28, \mathrm{~A} / \mathrm{A}_{\text {full }}=0.22
$$

- Calculate $\mathrm{A}_{\text {full }}$.

$$
\begin{aligned}
& A_{\text {full }}=\left(\frac{\pi}{4}\right) D^{2}=\left(\frac{\pi}{4}\right)(0.3 m)^{2} \\
& A_{\text {full }}=0.0707 \mathrm{~m}^{2}
\end{aligned}
$$

## Sewers Flowing Partly Full

- Substituting:

$$
\begin{aligned}
& \frac{A}{A_{\text {full }}}=0.22=\frac{A}{0.0707 \mathrm{~m}^{2}} \\
& A=0.0156 \mathrm{~m}^{2}
\end{aligned}
$$

- Substituting into the continuity equation:

$$
\begin{aligned}
& Q=V A \\
& 0.01 \mathrm{~m}^{3} / \mathrm{sec}=V\left(0.0156 \mathrm{~m}^{2}\right) \\
& V=0.641 \mathrm{~m} / \mathrm{sec}
\end{aligned}
$$

## In-Class Problem (Partially Flowing Sewer)

- Determine the depth of flow and velocity in a sewer with a diameter of 600 mm having a HGL slope of $0.005 \mathrm{~m} / \mathrm{m}$ with an n value of 0.013 when discharging $0.055 \mathrm{~m}^{3} / \mathrm{sec}$.


## Remember the problem having two "correct" answers:



Figure 7-7 Transition in an open channel.
The specific energy diagram is used to determine the most likely water depth.
Prasuhn 1987

## Typical Specific Energy Diagram

## (Figure 3.6, Chin 2006)

$t E=y+\alpha \frac{V^{2}}{2 g}$
If water depth is deeper than the
critical depth $\left(y_{c}\right)$, then the flow is

If the water depth is shallower than the critical depth, then the flow is supercritical. critical, small changes in specific energy results in large depth changes, resulting in unstable and excessive wave action. Fr should be $<0.86$ or $>1.13$ to prevent this.


Specific energy, $E$


Figure 7-8 Specific energy diagram.

Stake of Fiow m Open Channals, es deter mined
by Rayuolts and Fronke numbern


Reywaids number: ratio of inertin to vosests:



Froub number: ruto of inortial feces to grauts foras

$g=$ acceleotion of gauth ( $32,2 \mathrm{zs}$ /ise
= hyboutcic deth $\frac{A}{T}: T=$ wden ftre
of $F=1$; than entical fiow
if $F$ <1; than strecrituel fiow
is FD : then sperectition siow

In this example, the upstream flow is subcritical (deeper than the critical depth). When the water approaches an upward step, the specific energy calculation results in two possible water depth solutions. The correct downstream water depth solution is on the same limb of the specific energy downstream water depth solution is on the same limb of the specific energy
diagram as the upstream water depth. In this case, the flow is still subcritical, although the water depth actually decreases (but it cannot pass through the
critical depth value).

(a) Specific energy diagram

(b) Water surface profile

Figure 7-9 Subcritical transition with upward step.

In this example, the upstream flow is supercritical (shallower than the critical depth). When the water approaches an upward step, the specific energy calculation results in two possible water depth solutions. The correct downstream water depth solution is on the same limb of the specific energy diagram as the upstream water depth. In this case, the flow is still supercritical: the water depth increases (but it cannot pass through the critical depth value).

(a) Specific energy diagram

(b) Water surface profile

Figure 7-10 Supercritical transition with upward step.

## Example Problem: Determine the Downstream Water Depth when Affected by Channel Bottom Rise (ex. 7-3, Prasuhn 1987)

Determine the downstream depth in a horizontal rectangular channel in which the bottom rises 0.5 ft , if the steady flow discharge is 300 cfs , the channel width is 12 ft , and the upstream depth is 4 ft .

The discharge per unit width is:
$q=\frac{Q}{b}=\frac{300 \mathrm{ft}^{3} / \mathrm{sec}}{12 \mathrm{ft}}=25 \mathrm{cfs} / \mathrm{ft}$

The critical depth is therefore:
$y_{c}=\left(\frac{q^{2}}{g}\right)^{1 / 3}=\left(\frac{(25 c f s / f t)^{2}}{32.2 f t / \sec ^{2}}\right)^{1 / 3}=2.69 f t$

The upstream depth is therefore subcritical.

(a) Specific energy diagram

(b) Water surface profile

Figure 7-9 Subcritical transition with upward step.

From the subcritical position on the specific energy diagram, the depth and the water surface elevation will both decrease downstream over the "bump" in the channel bottom. Therefore, $\mathrm{y}_{2}$ must be greater than $\mathrm{y}_{\mathrm{c}}$ and less than $y_{1}-\Delta z$ :
$2.69 \mathrm{ft}<\mathrm{y}_{2}<3.5 \mathrm{ft}$
Solving the equation by iteration within this range results in the solution of $\mathrm{y}_{2}=3.09 \mathrm{ft}$. The trial solutions can be used to draw in the specific energy diagram.

The upstream specific energy is calculated to be:

$$
H_{01}=y_{1}+\frac{q^{2}}{2 g y_{1}^{2}}=4 f t+\frac{(25 c f s / f t)^{2}}{2\left(32.2 f t / \sec ^{2}\right)(4 f t)^{2}}=4.61 f t
$$

and the corresponding downstream specific energy is:

$$
\begin{aligned}
& H_{02}=H_{01}-\Delta z=y_{2}+\frac{q^{2}}{2 g y_{2}^{2}} \\
& H_{02}=4.61 f t-0.5 f t=4.11 f t=y_{2}+\frac{(25 c f s / f t)^{2}}{2\left(32.2 f t / \mathrm{sec}^{2}\right) y_{2}^{2}}
\end{aligned}
$$



## "Choke"

What happens when $\Delta z$ is increased to a greater and greater value under subcritical conditions? As $\Delta z$ increases, $\mathrm{H}_{02}$ must also continue to decrease. Therefore, $y_{2}$ decreases as well. The limit is reached for subcritical flow when $y_{2}$ equals the critical depth at which point the transition becomes a "choke." A further increase in $\Delta z$ results in the impossible situation where $\mathrm{H}_{02}$ is less than $\mathrm{H}_{\text {omin }}$ (there would be no positive solution to $\mathrm{H}_{\mathrm{o} 2}$ ): the upstream flow has insufficient energy to pass through the transition at the specified discharge.

The flow will not cease, but will adjust itself to either a lower discharge or an increase in specific energy. The flow will likely not change due to upstream flow sources. The upstream flow will increase both its upstream depth and specific energy by means of a gentle swell or a series of small waves that travel upstream. The new upstream depth will be such that the flow can just pass the transition and $y_{2}$ will equal $y_{c}$, and $H_{02}$ will equal $H_{o m i n}$. $H_{01}$ will exceed $H_{o m i n}$ by the value of $\Delta z$. This transition is called a choke since the critical depth prevails regardless of the increase in upstream energy

During supercritical flow conditions, the flow behavior is different as $\Delta z$ increases. A choke occurs when the minimum specific energy is reached. However, when additional $\Delta \mathrm{z}$ occurs, a surge (wall of water) moves upstream. When equilibrium is reached, the supercritical flow will have been replaced by the identical subcritical flow discussed above, and the transition will continue to act as a choke.
(summarized from Prasuhn 1987)

